# 470. The Electronic Orbitals, Shapes, and Spectra of Polyatomic Molecules. Part V.* Tetratomic, Non-hydride Molecules, $\mathrm{AB}_{3}$. 

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The procedure developed in previous papers of this series is applied to tetratomic, non-hydride molecules, $\mathrm{AB}_{3}$. A correlation diagram is plotted between the orbitals possible for planar $\mathrm{AB}_{3}$ molecules and those possible for pyramidal molecules. This enables a discussion to be given of the shapes of such molecules, and comments to be made on their reactivities and spectra. $\mathrm{AB}_{3}$ molecules containing $\ngtr 24$ valency electrons are planar in their ground states. Molecules containing 25 or 26 valency electrons should be pyramidal in their ground states, the BAB angle decreasing with increase in the number of electrons. Molecules containing 24 valency electrons should be pyramidal in their first excited states. The ground states of molecules containing 28 valency electrons should be planar or nearly planar.

Orbitals of Non-hydride, $\mathrm{AB}_{3}$ Molecules.-The lowest-lying possible orbitals for a planar $\mathrm{AB}_{3}$ molecule may be approximately described as follows. (i) Three $s$ lone-pair orbitals, one on each of the B atoms. In known $\mathrm{AB}_{3}$ molecules, B is usually of high electronegativity, greater than that of A . To a first approximation, therefore, these $s$ orbitals may be supposed so tightly bound that they hardly interact with each other or with other orbitals. They remain essentially atomic in character and vary little in binding energy as the BAB angle is varied. (ii) Three bonding orbitals. Regarded as localized orbitals, they are formed from $s p^{2}$ hybrid valencies of A overlapping in-phase with $p$ valencies of the B atoms that point towards atom A. Regarded as non-localized orbitals, they are of species $a_{1}{ }^{\prime}, e^{\prime}, e^{\prime}$, as described for $\mathrm{AH}_{3}$ molecules. (iii) Three orbitals which are approximately $p$ lone-pair orbitals, one on each of the B atoms. The axis of each of these $p$ orbitals lies in the molecular plane and perpendicular to the adjacent $\mathrm{B}-\mathrm{A}$ line. The orbitals are approximately non-bonding because the B atoms are far apart. They form

[^0]non-localized orbitals as indicated in Fig. 1. Two of the orbitals are degenerate and of species $e^{\prime}$. They have a net slight bonding effect between the $B$ atoms. The third orbital is of species $a_{2}{ }^{\prime}$ and is slightly anti-bonding in each $\mathrm{B}-\mathrm{B}$ distance; it therefore is of slightly higher energy than the $e^{\prime}$ orbitals. (iv) Four orbitals built from $p$ orbitals, on the $A$ and the three $B$ atoms, which have their axes perpendicular to the molecular plane (see Fig. 2). The lowest-energy orbital of these has all four $p$ orbitals overlapping inphase. It is of species $a_{2}{ }^{\prime \prime}$ and is bonding in each $\mathrm{A}-\mathrm{B}$ and $\mathrm{B}-\mathrm{B}$ distance. The next two orbitals each have a nodal plane perpendicular to the molecular plane and passing through atom A . They are degenerate, the two-fold degeneracy arising from the two independent and mutually perpendicular possible settings of the above nodal plane. Their species symbol is thus $e^{\prime \prime}$. They have a net anti-bonding effect between the B atoms and are $\mathrm{A} \leftrightarrow \mathrm{B}$ non-bonding. The fourth orbital has the $p$ orbital on atom A overlapping out-ofphase with each of the $p$ orbitals on the B atoms. It is of species $a_{2}{ }^{\prime \prime}$ and is anti-bonding in each A-B distance but bonding between the three B atoms. To denote its $\mathrm{A} \leftrightarrow \mathrm{B}$ antibonding character it is written ${\overrightarrow{a_{2}}}^{\prime \prime}$. The fact that B usually has a higher electronegativity than A means that the first three orbitals are more localized on the B atoms than on A , while the fourth orbital is more localized on atom A. The argument is the same as that used in Part II in the discussion of the localization of orbitals. The bonding or anti-bonding properties of the orbitals in (iii) and (iv) suggest that the order of decreasing binding energy is $a_{2}{ }^{\prime \prime} e^{\prime} e^{\prime} e^{\prime \prime} e^{\prime \prime} a_{2}{ }^{\prime}$, followed by $\bar{a}_{2}{ }^{\prime \prime}$. (v) Three orbitals which are built by

Fig. 1. (The choice of form for the e orbitals is not unique.)


Fig. :. (The forms of the $\mathrm{e}^{\prime \prime}$ orbitals are not unique.)


the out-of-phase overlap of an orbital on A with an appropriate group orbital compounded from the three $\mathrm{B} p$ orbitals whose axes lie along the $\mathrm{A}-\mathrm{B}$ lines. In the lowest of these (species $a_{1}{ }^{\prime}$ ) the three $\mathrm{B} p$ orbitals overlap in-phase with each other but out-of-phase with a central $s$ orbital on A . The orbital is anti-bonding in each $\mathrm{A}-\mathrm{B}$ distance and is therefore written $\vec{a}_{1}{ }^{\prime}$. It corresponds to the $\bar{a}_{1}{ }^{\prime}$ orbital described for $\mathrm{AH}_{3}$ molecules, and is more localized on A than on the B atoms.

The $\bar{a}_{2}{ }^{\prime \prime}$ orbital corresponds to the $a_{2}{ }^{\prime \prime}$ orbital described for $\mathrm{AH}_{3}$ molecules. Whereas the latter was non-bonding, however, the former is anti-bonding. In $\mathrm{AH}_{3}$ molecules the $a_{2}{ }^{\prime \prime}$ undoubtedly lies well below the $\vec{a}_{1}{ }^{\prime}$ orbital. In non-hydride $\mathrm{AB}_{3}$ molecules the antibonding character of $\bar{a}_{2}{ }^{\prime \prime}$ causes it to become of higher energy relative to the other orbitals. It is therefore uncertain in $\mathrm{AB}_{3}$ whether $\bar{a}_{2}^{\prime \prime}$ or $\bar{a}_{1}^{\prime}$ lies the higher. The point is exactly analogous to the question whether $\bar{\pi}_{u}$ or $\bar{\sigma}_{g}$ lies the lower for a linear $\mathrm{AB}_{2}$ molecule (Part II).

The lowest-lying orbitals for a pyramidal $\mathrm{AB}_{3}$ molecule in which the BAB angle is $90^{\circ}$ may be approximately described as follows. (i) Three $s$ lone-pair orbitals, one on each B atom. These remain essentially atomic in character and vary little in binding energy as the BAB angle is varied. In Fig. 3, which correlates the orbitals of pyramidal and planar
$\mathrm{AB}_{3}$ molecules,* they are therefore represented by three horizontal lines close together. (ii) Three bonding orbitals. Regarded as localized orbitals, these utilize pure $p$ valencies of A. Regarded as non-localized orbitals, they are of species $a_{1}, e, e$ and are represented by lines that fall from left to right in Fig. 3, as described for $\mathrm{AH}_{3}$ molecules. (iii) Six approximate lone-pair $p$ orbitals, two on each $B$ atom and each of axis perpendicular to the adjacent $A-B$ line. Since these orbitals are approximate lone-pair orbitals they are represented in Fig. 3 by lines that remain fairly close together whatever the angle BAB. Moreover, as the angle changes, they involve no changing hybridization of A valencies. Whether they increase or decrease in binding energy as the BAB angle increases is therefore determined by whether they are $\mathrm{B} \longleftrightarrow \mathrm{B}$ anti-bonding or bonding respectively [principle (iii) of Part I]. These $p$ orbitals may be separated into two groups; those whose axes lie in the plane of the three B atoms and those whose axes do not. Consider the former first. The non-localized combinations of the lone-pair orbitals are of species $e, e, a_{2}$, corresponding to the $e^{\prime}$, $e^{\prime}, a_{2}{ }^{\prime}$ group of the planar molecule. The $e, c$ orbitals are weakly bonding between the B atoms as a net effect. They therefore are more tightly bound when the B atoms are close together (angle BAB $=$ $90^{\circ}$ ) than when far apart (angle $\mathrm{BAB}=120^{\circ}$ ), and are represented in Fig. 3 by lines that rise from left to right. The $a_{2}$ orbital is anti-bonding between the B atoms. It is therefore more tightly bound in the planar molecule and is represented by a line that falls from left to right.

Considering now the remaining three $\mathrm{B} p$ orbitals: these give rise to $a_{1}, e, e$ orbitals corresponding to the $a_{2}{ }^{\prime \prime}, e^{\prime \prime}, e^{\prime \prime}$ group of the planar molecule. The $e, e$ or $e^{\prime \prime}, e^{\prime \prime}$ orbitals have a net anti-bonding effect between the B atoms. They are therefore represented in Fig. 3 by lines that fall from left to right. The $a_{1}$ or $a_{2}{ }^{\prime \prime}$ orbital is bonding between the B atoms and might therefore be expected to rise from left to right in Fig. 3. On the other hand, it probably changes from being built almost solely from B orbitals in the $90^{\circ}$ molecule to being built in-

Fig. 3.
 creasingly also from an A $p$ orbital as the BAB angle increases, i.e., it increases in $\mathrm{A} \longleftrightarrow \mathrm{B}$ bonding character. As a result, the orbital may vary comparatively little with change of angle and has been represented in the Figure by a horizontal line. In any case, however, the variation in the $a_{2}{ }^{\prime \prime}, e^{\prime \prime}, e^{\prime \prime}$ or $a_{2}{ }^{\prime}, e^{\prime}, e^{\prime}$ group will be small compared with the variation in the $a_{1}{ }^{\prime}, e^{\prime}, e^{\prime}$ group. (iv) An orbital that is approximately an $s$ orbital on atom A. It is of species $a_{1}$ (written $a_{1} s$ ). As the BAB angle increases, this orbital rises steeply in energy, as in $\mathrm{AH}_{3}$, because it becomes more and more built from a $p$ orbital of $A$. In $\mathrm{AB}_{3}$, in addition, it becomes more and more $\mathrm{A} \leftrightarrow \mathrm{B}$ antibonding, tending to approach the $\bar{a}_{2}^{\prime \prime}$ orbital of the planar molecule. (v) An orbital of species $a_{1}$ (written $\bar{a}_{1}$ ) which is anti-bonding in each A-B distance and corresponds to the $\bar{a}_{1}{ }^{\prime}$ orbital of the planar molecule. It is built from an $s$ orbital of A in the planar molecule, but from a $p_{2}$ orbital of A in the $90^{\circ}$ molecule and is therefore represented in Fig. 3 by a curve that falls steeply from left to right. If $\bar{a}_{1}$ for the planar molecule lies below $\bar{a}_{2}{ }^{\prime \prime}$, the steep curve from the $a_{1} s$ orbital will lead eventually to $\bar{a}_{1}{ }^{\prime}$. Otherwise, curves of the same species would cross. Similarly, the curve from the $\dot{a}_{1}{ }^{\prime}$ orbital would then lead not to the $\vec{a}_{1}{ }^{\prime}$ but to the $\vec{a}_{2}{ }^{\prime \prime}$ orbital.

Shapes of $\mathrm{AB}_{3}$ Molecules.-Of the twelve lowest orbital curves in Fig. 3, all but two either fall from left to right or remain roughly constant in energy. The two that rise only

[^1]do so slightly. One therefore has a strong expectation that $\mathrm{AB}_{3}$ molecules containing 24 or less valency electrons will be planar. This is borne out by the facts (see Table). A molecule with 25 or 26 valency electrons, however, has to have one or two electrons, respectively, in the $a_{1} S-\bar{a}_{2}^{\prime \prime}$ or $\bar{a}_{1}^{\prime}$ orbital. This means that such molecules can be expected to be pyramidal, the BAB angle being less for a 26 - than for a 25 -electron molecule. There appears to be only one example of a 25 -electron molecule, viz., $\mathrm{ClO}_{3}$ which, although probably pyramidal, has an unknown OClO angle. 26 -Electron molecules are all pyramidal as expected, the BAB angle being $\sim 100^{\circ}$. As with triatomic molecules, one notices the importance of pure number in determining the shape of a molecule. The Table includes a few $\mathrm{B}_{2} \mathrm{AC}$ molecules because a correlation diagram very similar to Fig. 3 (except that no orbitals are now degenerate) should apply to such molecules (see following papers). The linkage of shape with number of valency electrons should therefore be the same for $\mathrm{B}_{2} \mathrm{AC}$ as for $\mathrm{AB}_{3}$ molecules.

Shapes of tetratomic $\mathrm{AB}_{3}$ (and some $\mathrm{B}_{2} \mathrm{AC}$ ) molecules.

No. of valency electrons
24
$\mathrm{CO}_{3}{ }^{--}$
$\mathrm{SO}_{3}$
$\mathrm{BF}_{3}$
$\mathrm{BCl}_{3}$
$\mathrm{BBr}_{3}$
$\mathrm{COCl}_{2}$
$\mathrm{COBr}_{2}$ Molecule
 $\mathrm{ClO}_{3}$ 26
25

$$
\begin{gathered}
\\
\\
\mathrm{IO}_{3}- \\
\mathrm{ClO}_{3}
\end{gathered}
$$



Reactivity of $\mathrm{AB}_{3}$ Molecules.-Fig. 3 also enables one to discuss the reactivities of tetratomic molecules. A radical such as $\mathrm{NO}_{3}$, for example, should have its odd electron in an $e^{\prime \prime}$ orbital which, we have argued, is largely localized on the O atoms. We have already seen (Part II) that the odd electron of the $\mathrm{NO}_{2}$ molecule is localized largely on the N atom. One would expect therefore the structure of $\mathrm{N}_{2} \mathrm{O}_{5}$ (which is probably formed from $\mathrm{NO}_{3}+\mathrm{NO}_{2}$ ) in the gas phase to be $\underset{\mathrm{O}}{\mathrm{O}}-\mathrm{N} \cdot \mathrm{O} \cdot \mathrm{N}<\mathrm{O}_{\mathrm{O}}^{\mathrm{O}}$ rather than ${ }_{\mathrm{O}}^{\mathrm{O}}-\mathrm{N} \cdot \mathrm{N}_{-}^{\prime}-\mathrm{O}$. This is known to be correct (Sidgwick, op. cit.).* Similarly, one expects the $\mathrm{N}_{2} \mathrm{O}_{6}$ molecule to be ${ }_{\mathrm{O}}^{\mathrm{O}}-\mathrm{N} \cdot \mathrm{O} \cdot \mathrm{O} \cdot \mathrm{N}<\mathrm{O}_{\mathrm{O}}^{\mathrm{O}}$ rather than the ethane-like structure $\mathrm{O}_{\mathrm{O}}^{\mathrm{O}} \mathrm{N} \cdot \mathrm{N}_{-}^{\mathrm{O}}-$. This again appears to be true (Sidgwick, op. cit.). The odd electron of the NO molecule lies in an orbital more localized on the N than on the O atom (cf. Part II). It follows that the $\mathrm{N}_{2} \mathrm{O}_{3}$ molecule is expected to have the structure $\mathrm{O} \cdot \mathrm{N} \cdot \mathrm{N}<\mathrm{O}$. Similarly, $\mathrm{N}_{2} \mathrm{O}_{4}$ (which dissociates to $2 \mathrm{NO}_{2}$ rather than to $\mathrm{NO}+\mathrm{NO}_{3}$ ) should have the structure ${ }_{\mathrm{O}}^{\mathrm{O}}-\mathrm{N} \cdot \mathrm{N}-<_{\mathrm{O}}^{\mathrm{O}}$ rather than $\mathrm{O} \cdot \mathrm{N} \cdot \mathrm{O} \cdot \mathrm{N}<_{\mathrm{O}}^{\mathrm{O}}$ or $\mathrm{O} \cdot \mathrm{N} \cdot \mathrm{O} \cdot \mathrm{O} \cdot \mathrm{N} \cdot \mathrm{O}$. Pauling suggested the second of these because the first violated his "adjacent charge" rule. However, there is now little doubt that the first is correct (Ingold and Ingold, Nature, 1947, 159, 743; Walsh, J. Chem. Phys., 1947, 15, 688). Sidgwick (op. cit.) gives the structure ONONO for the $\mathrm{N}_{2} \mathrm{O}_{3}$ molecule, but the evidence for this does not appear unambiguous. NO should react with $\mathrm{NO}_{3}$ to give the molecule $\mathrm{O} \cdot \mathrm{N} \cdot \mathrm{O} \cdot \mathrm{N}<\mathrm{O}$ which could give $2 \mathrm{NO}_{2}$. The step $\mathrm{NO}+\mathrm{NO}_{3}=2 \mathrm{NO}_{2}$ may occur in the reaction between NO and $\mathrm{O}_{2}$ to give (finally) $\mathrm{NO}_{2}$.

Similar considerations should govern the interaction of $\mathrm{PO}, \mathrm{PO}_{2}$, and $\mathrm{PO}_{3}$ radicals, and are relevant to an understanding of how $\mathrm{P}_{4}$ on oxidation produces $\mathrm{P}_{4} \mathrm{O}_{10}$ molecules.

Spectra of $\mathrm{AB}_{3}$ Molecules.-By the Franck-Condon Principle, since the first excited states of such molecules as $\mathrm{CO}_{3}^{--}, \mathrm{NO}_{3}^{-}, \mathrm{BF}_{3}$, or $\mathrm{SO}_{3}$ should be non-planar, one expects the longest-wave-length electronic transitions of these molecules to occupy a wide region of the spectrum. This is true of, e.g., $\mathrm{SO}_{3}$, which has a long-wave-length transition extending from 3000 to beyond $2200 \AA$. A further reason lies in the increase of $\mathrm{A}-\mathrm{B}$ length that will occur in the transition. The $\bar{a}_{2}{ }^{\prime \prime}$ and the $\bar{a}_{1}{ }^{\prime}$ orbital are both so highly $\mathrm{A} \longleftrightarrow \mathrm{B}$ anti-bonding that it is not surprising that all the known spectra of 26 -electron molecules are continuous. The spectra of $\mathrm{PCl}_{3}, \mathrm{PBr}_{3}, \mathrm{AsCl}_{3}, \mathrm{SbCl}_{3}$, and $\mathrm{BiCl}_{3}$ (Sponer and Teller, Rev. Mod. Phys., 1941, 13, 75) supply examples.

The $\mathrm{SO}_{3}$ absorption referred to above consists of weak, diffuse, bands superimposed on a continuous background. We suggest that the bands and the continuous background are to be interpreted as due to two transitions which may both be formulated

$$
\begin{equation*}
\left.\cdots\left(a_{2}\right)\left(a_{1}\right),{ }^{1} A_{2} \longleftarrow \cdots\left(a_{2}\right)^{\prime}\right)^{2}{ }^{1} A_{1}^{\prime} \tag{1}
\end{equation*}
$$

corresponding to the two transitions

$$
\begin{align*}
& \cdots\left(a_{2}{ }^{\prime}\right)\left(\bar{a}^{\prime \prime}\right),{ }^{1} A_{1}{ }^{\prime \prime} \longleftarrow \cdots\left(a_{2}{ }^{\prime}\right)^{2},{ }^{1} A_{1}{ }^{\prime}  \tag{2}\\
& \left.\cdots\left(a_{2}\right)^{\prime}\left(\overline{a_{1}}\right),,^{1} A_{2}{ }^{\prime} \longleftarrow \cdots\left(a_{2}\right)^{2}\right)^{2},{ }^{1} A_{1}
\end{align*}
$$

and
Both (2) and (3) are forbidden; (2) should lead to a pyramidal and (3) to a planar upper state. In agreement the absorption is not very strong, $\varepsilon$ slowly increasing from a very low value at $3000 \AA$ to about 100 at $2150 \AA$ (Fajans and Goodeve, Trans. Faraday Soc., 1936, 32, 511). It should be noted, however, that the peak of the absorption lies at some unknown wave-length below $2150 \AA$. The banded part of the absorption appears to begin at about $2700 \AA$, but the bands have only been measured below $2500 \AA$. They are separated by an approximately constant frequency difference of
${ }^{*}$ In the solid state, $\mathrm{N}_{2} \mathrm{O}_{5}$ is known to be $\mathrm{NO}_{2}{ }^{+}, \mathrm{NO}_{3}^{-}$(Grison, Eriks, and de Vries, Acta Cryst., 1950, 3, 290).
$430 \mathrm{~cm} .^{-1}$. The probable interpretation of this is that it represents either the symmetrical stretching or the symmetrical bending vibration (both of species $a_{1}$ for a pyramidal upper state, or $a_{1}{ }^{\prime}$ and $a_{2}{ }^{\prime \prime}$ respectively for a planar form) of the upper state. In the ground state the $a_{1}{ }^{\prime}$ - and the $a_{2}{ }^{\prime \prime}$-vibration have frequencies of 1069 and $653 \mathrm{~cm} .^{-1}$ respectively (Gerding and Lecomte, Physica, 1939, 6, 737). The bands have a curious flat-topped structure (see Figure given by Fajans and Goodeve, loc. cit.). Indeed, $430 \mathrm{~cm} .^{-1}$ really refers to the measured separations of the minima (which alone are sharp) between the bands. No explanation of this band shape has yet been given.

The 23 -electron $\mathrm{NO}_{3}$ molecule is known to give rise to absorption bands between 7100 and $5000 \AA$ (see Jones and Wulf, J. Chem. Phys., 1937, 5, 873). The extinction coefficient is quite high ( $\sim 2000$ at $\lambda_{\text {max }}$. which is $\sim 6650 \AA$ ). This suggests the identification (see Fig. 3)

$$
\begin{equation*}
\left.\cdots\left(e^{\prime}\right)^{3}\left(a_{2}\right)^{\prime}\right)^{1} E^{\prime} \longleftarrow \cdots\left(e^{\prime}\right)^{4}\left(a_{2}{ }^{\prime}\right),{ }^{1} A_{2}{ }^{\prime} \tag{4}
\end{equation*}
$$

Now (4) is an allowed transition, the molecule being planar in the equilibrium forms of both combining states. Not surprisingly, therefore, the bands are quite different from those of $\mathrm{SO}_{3}$, being comparatively sharp and narrow. Moreover, the very long wave-length of this $\mathrm{NO}_{3}$ absorption system is in accord with the discussion above and with the closeness of the binding energies of the orbitals $a_{2}{ }^{\prime \prime}, e^{\prime \prime}, e^{\prime \prime}$ and $e^{\prime}, e^{\prime}, a_{2}{ }^{\prime}$ in Fig. 3. If (4) is assumed to be correct, the bands should turn out under higher dispersion to be perpendicular in type. Since the ( $e^{\prime}$ ) and ( $a_{2}{ }^{\prime}$ ) orbitals concerned are respectively $\mathrm{O} \longleftrightarrow 0$ bonding and anti-bonding, the main vibrational frequency involved should be the $a_{1}{ }^{\prime}$ totally symmetrical stretching vibration of the upper state.

The absorption shown by dilute aqueous solutions of nitrates (a rounded peak between 3000 and $2500 \AA$ with $\varepsilon \sim 7$; Sponer and Teller, loc. cit.; Katzin, J. Chem. Phys., 1950, 18, 789) is probably to be interpreted as transition (1) of the nitrate ion, more or less perturbed by the adjacent solvent and cations.

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[^0]:    * Part IV, preceding paper.

[^1]:    * For clarity, Fig. 3 (like Fig. I of Part IV) is over-simplified. Thus the $a_{1} s_{\Delta}$ and $e e$ bonding orbital curves probably cross (see Part VII). Crossings (or avoided crossings) would complicate the diagram, but need not upset the basic ideas of this paper.

